

9. ELECTROMAGNETIC WAVES

9.1 Classical wave equation

We saw in PHYS1231 (!) how Newton's law $\mathbf{F} = m\mathbf{a}$ can be applied to an element of **string** under tension to derive the differential equation describing waves propagating with velocity v .

$$\frac{\partial^2 f}{\partial z^2} = \frac{1}{v^2} \frac{\partial^2 f}{\partial t^2} \quad \text{classical wave equation}$$

(in the case of waves on a string $v = \sqrt{T/\mu}$; T = tension, μ = string's mass per unit length).

(see Tipler p446-448 (wave on string) and p1012-1014 (EM wave))

9.2 Wave equation solutions

Solutions of the classical wave equation are functions with the form

$$f(z, t) = g(z \pm vt)$$



+ describes a wave travelling in -ve z -dirⁿ.
- describes a wave travelling in +ve z -dirⁿ.

and since the wave equation is *linear* (there are no differentials raised to power 2 or higher) we can have superpositions of solutions too:

$$f(z,t) = \underbrace{g(z - vt) + h(z + vt)}$$

wave travelling toward +ve z plus wave travelling toward -ve z is also a solution of the wave equation

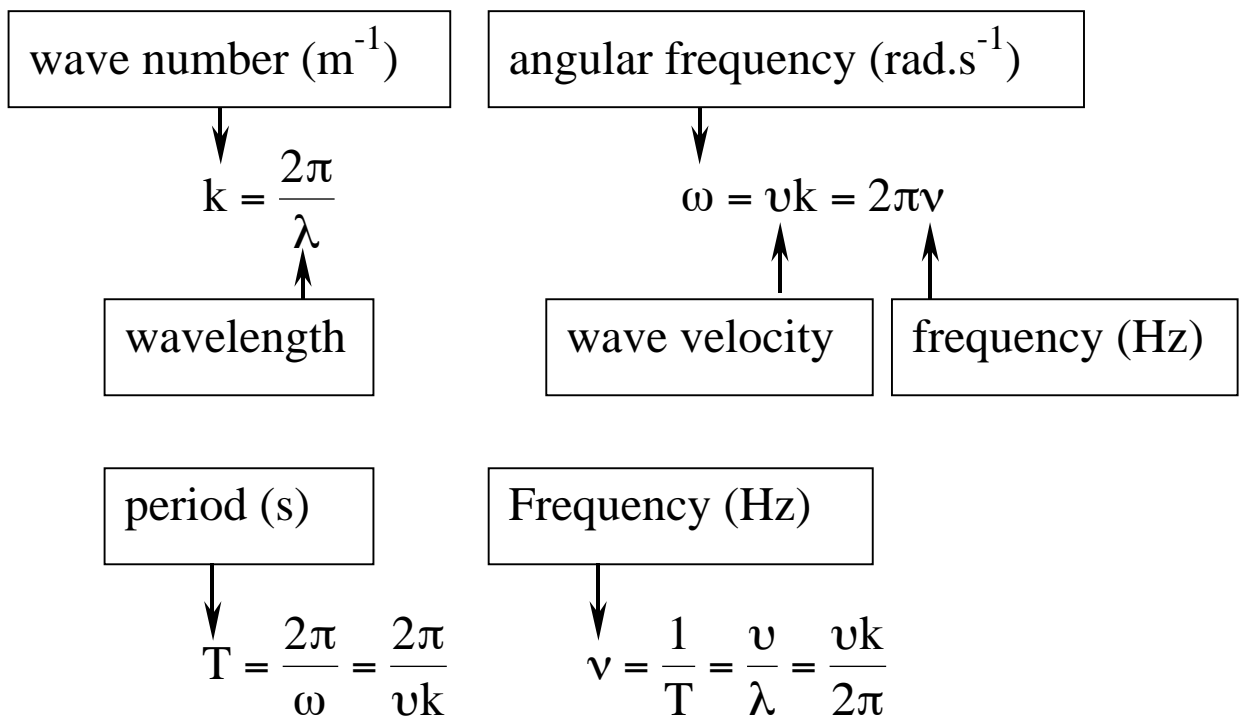
Of all the (many!) possible mathematical functions g that are solutions to the wave equation, harmonic (sine and cosine) functions are the most useful:

$$f(z,t) = A \cos[k(z - vt) + \delta]$$

Diagram illustrating the components of the harmonic wave function $f(z,t) = A \cos[k(z - vt) + \delta]$:

- amplitude** points to A .
- wave number** points to k .
- velocity** points to v .
- phase constant** points to δ .
- phase** (bracketed) points to the entire argument $k(z - vt) + \delta$.

We can re-write $f(z,t)$ in various useful forms using



A very useful form of solution is

$$f(z,t) = A \cos(kz \pm \omega t \pm \delta)$$

9.3 Wave equation solution in complex notation

Recalling that $e^{i\theta} = \cos\theta + i\sin\theta$, we can write

$$f(z,t) = \text{Re} \left[A e^{i(kz - \omega t + \delta)} \right]$$

Re: take the real part of
this complex exponential

If we make the wave function complex by including a complex amplitude $\tilde{A} = Ae^{i\delta}$ we write:

$$\tilde{f}(z,t) = \tilde{A}e^{i(kz-\omega t)}$$

9.4 Maxwell's equations give propagating EM waves

If we take the two Maxwell eqns. (in differential form)

$$(iii) \quad \nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \quad (\text{Faraday's law})$$

$$(iv) \quad \nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \quad (\text{Ampère's law with Maxwell's correction})$$

and take the curl of these two in a region of space with

- *no charge*
- *no current*

Take curl of (iii):

$$\nabla \times (\nabla \times \mathbf{E}) = \nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E} = \nabla \times \left(-\frac{\partial \mathbf{B}}{\partial t} \right)$$

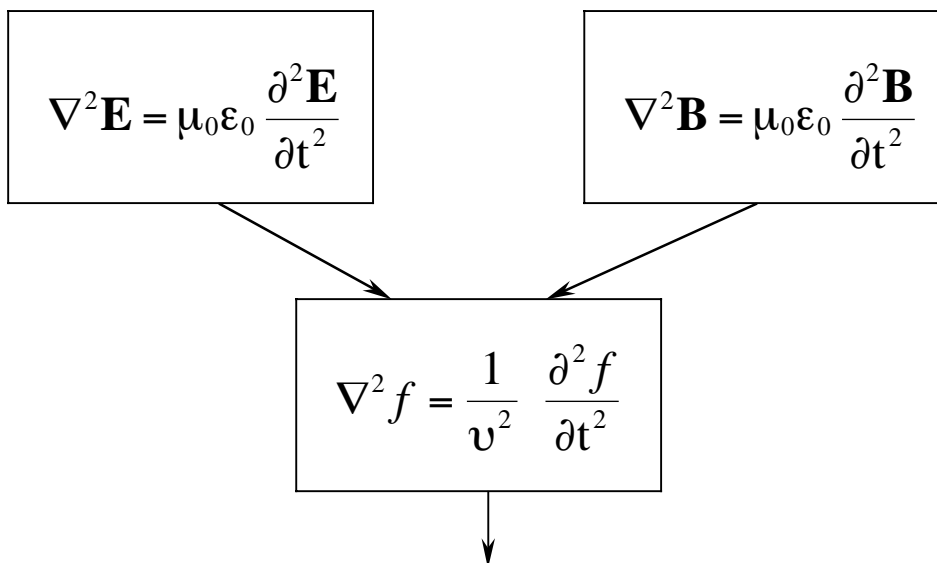
$$= -\frac{\partial}{\partial t}(\nabla \times \mathbf{B}) = -\mu_0 \epsilon_0 \frac{\partial^2 \mathbf{E}}{\partial t^2}$$

take curl of (iv)

$$\nabla \times (\nabla \times \mathbf{B}) = \nabla(\nabla \cdot \mathbf{B}) - \nabla^2 \mathbf{B} = \nabla \times \left(\mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \right)$$

$$= -\mu_0 \epsilon_0 \frac{\partial}{\partial t}(\nabla \times \mathbf{E}) = -\mu_0 \epsilon_0 \frac{\partial^2 \mathbf{B}}{\partial t^2}$$

The equations for \mathbf{E} and \mathbf{B}



with

$$v = \frac{1}{\sqrt{\epsilon_0 \mu_0}} = 3.00 \times 10^8 \text{ m/s}$$

These equations (obviously (!)) describe electromagnetic waves EM propagating at velocity $v = c$.

James Clerk Maxwell (1831-1879), Scottish mathematician and physicist in “*A Treatise on Electricity and Magnetism*” (1873) first showed that the (eponymous; Maxwell) equations implicitly require the existence of EM waves travelling at the speed of light.

Maxwell also knew the numerical values of μ_0 and ϵ_0 measured in Germany in 1856. The speed of light was also known.

History: Speed of light first investigated by Ole Römer in 1676. Römer observed the eclipse of Io, one of Jupiter’s moons: Römer found the speed of light to be very large...but finite!

Armand Fizeau used the Fizeau Wheel (a rotating toothed wheel and distant mirror) and found $c = 3.15 \times 10^8 \text{ ms}^{-1}$ in ~1850.

(<http://scienceworld.wolfram.com/physics/FizeauWheel.html> and links therein is an interesting source)

9.5 Power in EM waves

Electromagnetic waves transmit *information* and *power*. We are already familiar with the energy density (per unit volume) in static magnetic and electric fields:

$$u = \frac{1}{2} \left(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right)$$

and since the electric and magnetic components contribute equally (see Griffiths p 378, Example 9.2),

$$B^2 = \frac{1}{c^2} E^2 = \mu_0 \epsilon_0 E^2$$

and,

$$u = \epsilon_0 E^2 = \epsilon_0 E_0^2 \cos^2(kz - \omega t + \delta) \quad *$$


↑
energy density in the
electromagnetic field

The **Poynting vector S** gives the energy per unit area per second and is defined

$$\mathbf{S} = \frac{1}{\mu_0} (\mathbf{E} \times \mathbf{B})$$

and using * above,

$$\mathbf{S} = c\epsilon_0 E_0^2 \cos^2(kz - \omega t + \delta) \hat{\mathbf{z}} = cu \hat{\mathbf{z}}$$


 energy density x
 velocity of EM waves

Note that EM waves also carry momentum $\rho = \frac{1}{c^2} \mathbf{S}$ you can read about this on p381 of Griffiths)

For the sine squared and cos squared functions, averaged over a cycle,

$$\langle \sin^2 \rangle = \langle \cos^2 \rangle = 1/2$$

$$\left\{ \frac{1}{T} \int_0^T \cos^2(kz - 2\pi t/T + \delta) dt = 1/2 \right\}$$

so that $\langle u \rangle = \frac{1}{2} \epsilon_0 E_0^2$

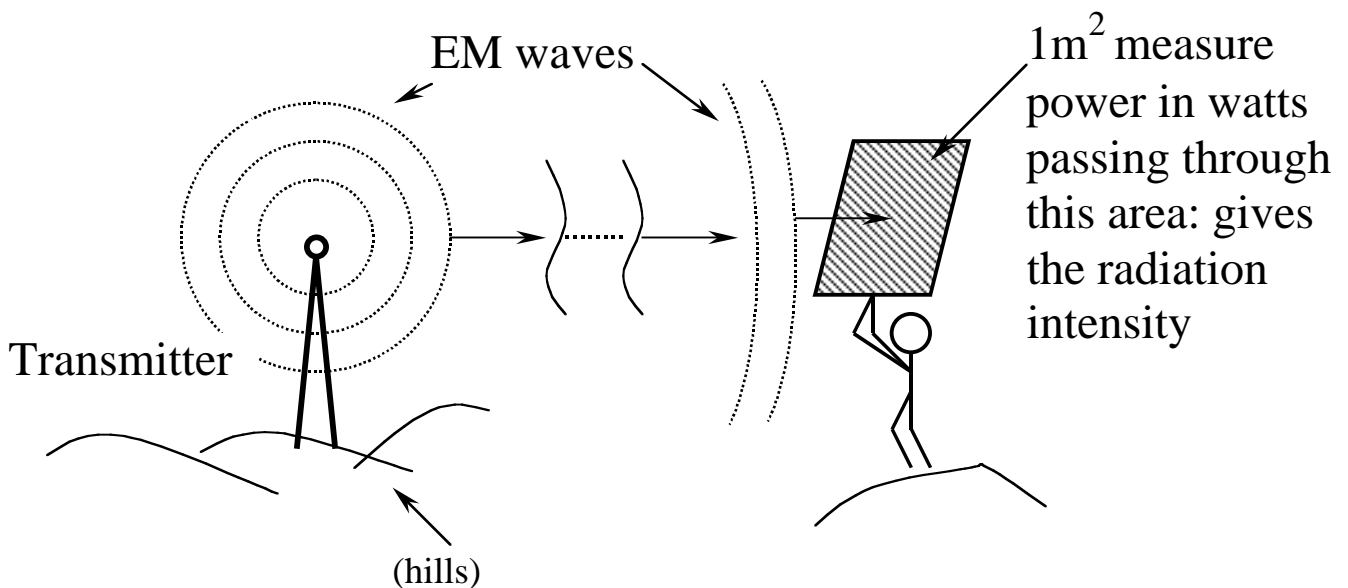
and

$$\langle \mathbf{S} \rangle = \frac{1}{2} c \epsilon_0 E_0^2 \hat{\mathbf{z}}$$

The EM wave intensity is

$$I \equiv \langle \mathbf{S} \rangle = \frac{1}{2} c \epsilon_0 E_0^2$$

This is the **average power per unit area** transmitted or transported by the wave. This is the intensity we met in PHYS1231 when discussing waves, in particular interference and diffraction of EM waves.



Problems 8, 9 and 10 on Problem Sheet 6 (the last sheet!) concern EM wave calculations.